

Bioengineering 280A
Principles of Biomedical Imaging

Fall Quarter 2009
MRI Lecture 2a

TT Liu, BE280A, UCSD Fall 2009

Gradient Fields

Define

$$\vec{G} \equiv G_x \hat{i} + G_y \hat{j} + G_z \hat{k} \quad \vec{r} \equiv x \hat{i} + y \hat{j} + z \hat{k}$$

So that

$$G_x x + G_y y + G_z z = \vec{G} \cdot \vec{r}$$

Also, let the gradient fields be a function of time. Then the z-directed magnetic field at each point in the volume is given by :

$$B_z(\vec{r}, t) = B_0 + \vec{G}(t) \cdot \vec{r}$$

TT Liu, BE280A, UCSD Fall 2009

Static Gradient Fields

In a uniform magnetic field, the transverse magnetization is given by:

$$M(t) = M(0) e^{-j\omega_0 t} e^{-t/T_2}$$

In the presence of non time-varying gradients we have

$$\begin{aligned} M(\vec{r}) &= M(\vec{r}, 0) e^{-j\gamma B_z(\vec{r}) t} e^{-t/T_2(\vec{r})} \\ &= M(\vec{r}, 0) e^{-j\gamma (B_0 + \vec{G} \cdot \vec{r}) t} e^{-t/T_2(\vec{r})} \\ &= M(\vec{r}, 0) e^{-j\omega_0 t} e^{-j\gamma \vec{G} \cdot \vec{r} t} e^{-t/T_2(\vec{r})} \end{aligned}$$

TT Liu, BE280A, UCSD Fall 2009

Time-Varying Gradient Fields

In the presence of time-varying gradients the frequency as a function of space and time is:

$$\begin{aligned} \omega(\vec{r}, t) &= \gamma B_z(\vec{r}, t) \\ &= \gamma B_0 + \gamma \vec{G}(t) \cdot \vec{r} \\ &= \omega_0 + \Delta\omega(\vec{r}, t) \end{aligned}$$

TT Liu, BE280A, UCSD Fall 2009

Phase

Phase = angle of the magnetization phasor
 Frequency = rate of change of angle (e.g. radians/sec)
 Phase = time integral of frequency

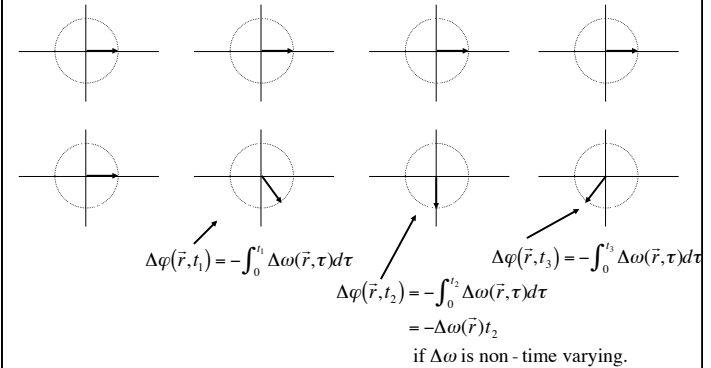
$$\begin{aligned}\varphi(\vec{r}, t) &= -\int_0^t \omega(\vec{r}, \tau) d\tau \\ &= -\omega_0 t + \Delta\varphi(\vec{r}, t)\end{aligned}$$

Where the incremental phase due to the gradients is

$$\begin{aligned}\Delta\varphi(\vec{r}, t) &= -\int_0^t \Delta\omega(\vec{r}, \tau) d\tau \\ &= -\int_0^t \gamma \vec{G}(\tau) \cdot \vec{r} d\tau\end{aligned}$$

TT Liu, BE280A, UCSD Fall 2009

Phase with constant gradient



TT Liu, BE280A, UCSD Fall 2009

Time-Varying Gradient Fields

The transverse magnetization is then given by

$$\begin{aligned}M(\vec{r}, t) &= M(\vec{r}, 0) e^{-t/T_2(\vec{r})} e^{i\varphi(\vec{r}, t)} \\ &= M(\vec{r}, 0) e^{-t/T_2(\vec{r})} e^{-j\omega_0 t} \exp\left(-j \int_0^t \Delta\omega(\vec{r}, \tau) d\tau\right) \\ &= M(\vec{r}, 0) e^{-t/T_2(\vec{r})} e^{-j\omega_0 t} \exp\left(-j\gamma \int_0^t \vec{G}(\tau) \cdot \vec{r} d\tau\right)\end{aligned}$$

TT Liu, BE280A, UCSD Fall 2009

Signal Equation

Signal from a volume

$$\begin{aligned}s_r(t) &= \int_V M(\vec{r}, t) dV \\ &= \int_x \int_y \int_z M(x, y, z, 0) e^{-t/T_2(\vec{r})} e^{-j\omega_0 t} \exp\left(-j\gamma \int_0^t \vec{G}(\tau) \cdot \vec{r} d\tau\right) dx dy dz\end{aligned}$$

For now, consider signal from a slice along z and drop the T_2 term. Define $m(x, y) = \int_{z_0 - \Delta z/2}^{z_0 + \Delta z/2} M(\vec{r}, t) dz$

To obtain

$$s_r(t) = \int_x \int_y m(x, y) e^{-j\omega_0 t} \exp\left(-j\gamma \int_0^t \vec{G}(\tau) \cdot \vec{r} d\tau\right) dx dy$$

TT Liu, BE280A, UCSD Fall 2009

Signal Equation

Demodulate the signal to obtain

$$\begin{aligned}
 s(t) &= e^{j\omega_0 t} S_r(t) \\
 &= \int_x \int_y m(x, y) \exp\left(-j\gamma \int_0^t \vec{G}(\tau) \cdot \vec{r} d\tau\right) dx dy \\
 &= \int_x \int_y m(x, y) \exp\left(-j\gamma \int_0^t [G_x(\tau)x + G_y(\tau)y] d\tau\right) dx dy \\
 &= \int_x \int_y m(x, y) \exp\left(-j2\pi(k_x(t)x + k_y(t)y)\right) dx dy
 \end{aligned}$$

Where

$$\begin{aligned}
 k_x(t) &= \frac{\gamma}{2\pi} \int_0^t G_x(\tau) d\tau \\
 k_y(t) &= \frac{\gamma}{2\pi} \int_0^t G_y(\tau) d\tau
 \end{aligned}$$

TT Liu, BE280A, UCSD Fall 2009

MR signal is Fourier Transform

$$\begin{aligned}
 s(t) &= \int_x \int_y m(x, y) \exp\left(-j2\pi(k_x(t)x + k_y(t)y)\right) dx dy \\
 &= M(k_x(t), k_y(t)) \\
 &= F[m(x, y)]_{k_x(t), k_y(t)}
 \end{aligned}$$

TT Liu, BE280A, UCSD Fall 2009

Recap

- Frequency = rate of change of phase.
- Higher magnetic field -> higher Larmor frequency -> phase changes more rapidly with time.
- With a constant gradient G_x , spins at different x locations precess at different frequencies -> spins at greater x -values change phase more rapidly.
- With a constant gradient, distribution of phases across x locations changes with time. (phase modulation)
- More rapid change of phase with x -> higher spatial frequency k_x

TT Liu, BE280A, UCSD Fall 2009

K-space

At each point in time, the received signal is the Fourier transform of the object

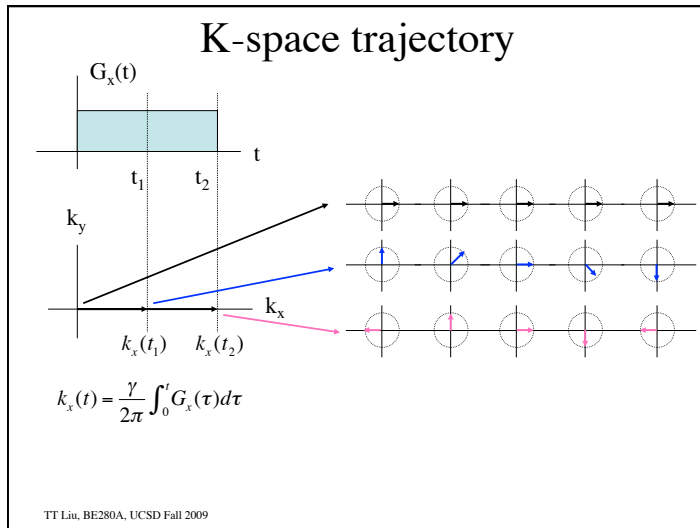
$$s(t) = M(k_x(t), k_y(t)) = F[m(x, y)]_{k_x(t), k_y(t)}$$

evaluated at the spatial frequencies:

$$\begin{aligned}
 k_x(t) &= \frac{\gamma}{2\pi} \int_0^t G_x(\tau) d\tau \\
 k_y(t) &= \frac{\gamma}{2\pi} \int_0^t G_y(\tau) d\tau
 \end{aligned}$$

Thus, the gradients control our position in k-space. The design of an MRI pulse sequence requires us to efficiently cover enough of k-space to form our image.

TT Liu, BE280A, UCSD Fall 2009



Units

Spatial frequencies (k_x, k_y) have units of 1/distance.
Most commonly, 1/cm

Gradient strengths have units of (magnetic field)/
distance. Most commonly G/cm or mT/m

$\gamma/(2\pi)$ has units of Hz/G or Hz/Tesla.

$$k_x(t) = \frac{\gamma}{2\pi} \int_0^t G_x(\tau) d\tau$$

$$= [\text{Hz}/\text{Gauss}][\text{Gauss}/\text{cm}][\text{sec}]$$

$$= [1/\text{cm}]$$

TT Liu, BE280A, UCSD Fall 2009

